

INTRODUCTION

Motivation: Estimate transport coefficients (mobility, thermal conductivity, shear viscosity).

Framework: Stochastic dynamics out of equilibrium.

Core idea: Make the flux the independent state variable instead of the forcing.

Hope: Derive estimators with better statistical properties (lower asymptotic variance).

GENERAL DEFINITION

Usual NEMD or "Thévenin" dynamics:

 $dX_t^{\eta} = b(X_t^{\eta}) dt + \sigma(X_t^{\eta}) dW_t + \eta F(X_t^{\eta}) dt,$

Reference (equilibrium) dynamics: $\eta = 0$. Norton dynamics:

$$\begin{cases} dY_t^r = b(Y_t^r) dt + \sigma(Y_t^r) dW_t + F(Y_t^r) \\ R(Y_t^r) = R(Y_0^r) = r. \end{cases}$$

Constrain the response flux R to be constant: dynamics on a submanifold of phase space.

Forcing process:

$$\mathrm{d}\Lambda^r_t = \lambda^r_t \,\mathrm{d}t + \mathrm{d}\widetilde{\Lambda}^r_t$$

with Λ^r a martingale.

TRANSPORT COEFFICIENTS

Thévenin:

$$\rho_{F,R} = \lim_{\eta \to 0} \frac{\mathbb{E}_{\eta} \left[R \right]}{\eta}$$

where \mathbb{E}_n is expectation with respect to steady-state of Thévenin dynamics.

Norton

$$\widetilde{\rho}_{F,R} = \lim_{r \to 0} \frac{r}{\mathbb{E}_r^* \left[\lambda\right]},$$

where \mathbb{E}_r^* is expectation with respect to steady-state of Norton dynamics, λ is the forcing.

In practice, these quantities are estimated through trajectory averages.



STOCHASTIC NORTON DYNAMICS Noé Blassel, Gabriel Stoltz

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 $) \mathrm{d}\Lambda_t^r,$

CLOSED-FORM

The dynamics can be written

$$\mathrm{d}Y_t^r = \overline{P}_{F,\nabla R}(Y_t^r) \left[b(Y_t^r) \mathrm{d}t + \sigma(Y_t^r) \mathrm{d}W_t \right] - \frac{\left(\nabla^2 R : \Pi_{F,\nabla R,\sigma}\right) (Y_t^r)}{2\nabla R(Y_t^r) \cdot F(Y_t^r)} F(Y_t^r) \,\mathrm{d}t,$$

with $\overline{P}_{F,\nabla R}(y)$ a non-orthogonal projector, and

$$\Pi_{F,\nabla R,\sigma}(y) = \left[\overline{P}_{F,\nabla R}\sigma\sigma^{\mathsf{T}}\overline{P}_{F,\nabla R}^{\mathsf{T}}\right](y).$$

orcing: $\lambda_t = \lambda(Y_t^r)$, with

Non-martingale part of the for

$$\lambda(y) = \left[-\frac{1}{F \cdot \nabla R} \left(b \cdot \nabla R + \frac{1}{2} \nabla^2 R : \Pi_{F, \nabla R, \sigma} \right) \right] (y).$$

Controllability condition: $F \cdot \nabla R \neq 0$.

Generalizations: Time-dependent flux $R(Y_t^r) = r_t$, multiple flux constraints $\mathbf{R}(Y_t^r) = \mathbf{r} \in \mathbb{R}^c$, also explicit.

LANGEVIN SETTING

dynamics:

$$\begin{cases} \mathrm{d}q_t = M^{-1}p_t \,\mathrm{d}t, \\ \mathrm{d}p_t = \overline{P}_{F,G}(q_t) \left(-\nabla V(q_t) \,\mathrm{d}t - \gamma M^{-1}p_t \,\mathrm{d}t + \sqrt{\frac{2\gamma}{\beta}} \,\mathrm{d}W_t \right) + \frac{\nabla G(q_t)p_t \cdot M^{-1}p_t}{F(q_t) \cdot G(q_t)} F(q_t) \,\mathrm{d}t \end{cases}$$

Mobility: $F \in \mathbb{R}^{dN}, G = M^{-1}F$. **Shear-viscosity:**

$$\forall 1 \le i \le N, \, \forall 2 \le \alpha \le d, \, F(q)_{i1} = f_y(q_{i2}), \, G(q)_{i1} = \frac{1}{m} \exp\left(\frac{2i\pi q_{i2}}{L}\right), \, F(q)_{i\alpha} = G(q)_{i\alpha} = 0$$

This corresponds to a longitudinal force with a transverse intensity profile. The response is the Fourier coefficient of the transverse profile for the longitudinal velocity.

Equivalence of responses



Take R(q, p) of the form $G(q) \cdot p$, adapted for mobility and shear viscosity computations. Norton

The Thévenin and Norton responses coïncide far outside of the linear response regime, and give the same transport coefficient. Computations were performed on a Lennard-Jones fluid of 1000 particles using a dedicated flux-preserving splitting implemented using the Julia package Molly.



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Norton estimators systematically outperform their Thévenin counterparts.





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ASYMPTOTIC VARIANCES





forcing

FUTURE QUESTIONS

- Existence and ergodicity results for the solutions to the Norton dynamics.
- Equivalence of Thévenin and Norton ensembles.
- Linear response theory/ Green-Kubo relations for the Norton ensemble.

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